

# Triangle Qubit Channel

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Given a triangle on a unit circle, we proved that an unital qubit channel can be constructed by cosine correlation functions. Furthermore, we proved that Bell's inequality is non-violation for the antipodal point of the triangle qubit channel. Finally, we discussed a triangle qubit channel formed by three cosine wave functions with relatively prime frequencies.

## I. INTRODUCTION

Quantum channels can transmit quantum states and classical information. In particular, a qubit channel is a quantum channel with a single qubit. Quantum channels play an important role in quantum computing, quantum communication, and quantum cryptography. In physics, we can implement quantum channels by the transmission of entangled photons through fiber optics or free space.

This paper originated from the author's previous research on Jones polynomials in quantum computing<sup>1</sup> and on Bell's inequality<sup>2</sup> of polynomial matrix<sup>3</sup>. In theorem 1, we constructed a qubit channel based on a triangle on the unit circle and figured out the cosine correlation functions as its elements.

To theorem 2, from K. Martin<sup>4</sup>'s research on the scope of quantum channels and the research of A. Fujiwara and P. Algoet<sup>5</sup> on Fujiwara-Algoet Condition (FAC), we proved the non-violation of Bell's inequality<sup>2</sup> for the antipodal point of an unital qubit channel. We also referred to the research of D. Braun<sup>6</sup> and colleagues on universal features of the quantum channels included the M.-D. Choi's matrix<sup>7</sup> for completely positive and trace-preserving.

## II. BACKGROUND

Define  $\mathcal{M}_d(\mathbb{C})$  to be the set of  $d \times d$  matrices over the complex field  $\mathbb{C}$  and  $\mathcal{D}_d(\mathbb{C})$  to be the set of  $d \times d$  density operator matrices that is positive, Hermitian, and trace one. A channel  $\phi : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$  is a linear map that is completely positive with trace-preserving. A channel  $\phi$  is also a map  $\phi : \mathcal{D}_d(\mathbb{C}) \rightarrow \mathcal{D}_d(\mathbb{C})$ . A unitary channel  $\phi_u : \mathcal{M}_d(\mathbb{C}) \rightarrow \mathcal{M}_d(\mathbb{C})$  is the set  $\phi_u(\rho) = U\rho U^\dagger$  where  $U$  is a unitary  $d \times d$  matrix and the operator  $\rho \in \mathcal{D}_d(\mathbb{C})$  is a density operator matrix.<sup>6</sup>

A qubit channel  $\phi : \mathcal{M}_2(\mathbb{C}) \rightarrow \mathcal{M}_2(\mathbb{C})$  is a two-dimensional channel. We can represent any state  $\rho \in \mathcal{M}_2(\mathbb{C})$  by Pauli matrices as the basis such that  $\rho = \frac{1}{2} \sum_{i=0}^3 r_i \sigma_i$  where  $r_i \in \mathbb{R}$  with  $r_0 = 1$ . The  $\text{trace}(\rho) = 1$  and  $r = (r_1, r_2, r_3)$  is the Bloch vector. A qubit channel  $\phi$  acting on state  $\rho \in \mathcal{M}_2(\mathbb{C})$  is a 4 by 4 real homogeneous matrix  $T_\phi$ . The positivity of Choi's matrix of  $T_\phi$  is equivalent to the complete positivity of the qubit channel  $\phi$ .

An unital qubit channel is a qubit channel  $\phi$  such that  $\phi(I/2) = I/2$ , that is the matrix  $T_\phi = \text{diag}(1, \lambda_1, \lambda_2, \lambda_3)$ . An antipodal point or an antipode of a point in a Bloch sphere is the diametrically opposite point. The antipodal map  $\phi' : \mathfrak{B}_3(\mathbb{R}) \rightarrow \mathfrak{B}_3(\mathbb{R})$  in Bloch sphere is  $\phi'(\lambda) = -\lambda$ .

## III. TRIANGLE QUBIT CHANNEL

**Theorem 1.** *Given a triangle on the unit circle, there exists an unital qubit channel described by the triple  $(\lambda_1, \lambda_2, \lambda_3) = (\cos(2\pi\omega_1)\cos(2\pi\omega_3), \cos(2\pi\omega_2)\cos(2\pi\omega_3), \cos(2\pi\omega_1)\cos(2\pi\omega_2))$  where the  $\omega_1, \omega_2, \omega_3 \in \mathbb{R}$ .*

*Proof.* Three vertices on the unit circle of the complex plane make a unique triangle. Given these three distinct vertices as a triple  $(\exp(i2\pi\theta_1), \exp(i2\pi\theta_2), \exp(i2\pi\theta_3))$  where  $\theta_1, \theta_2, \theta_3 \in \mathbb{R}$ , we define the triangle qubit matrix  $T(\Theta; K_4)$  with the basis of the Klein four-group  $K_4$  below,

$$T(\Theta; K_4) = \frac{1}{2} (\kappa_0 + \kappa_1 \exp(i2\pi\theta_1) + \kappa_2 \exp(i2\pi\theta_2) + \kappa_3 \exp(i2\pi\theta_3)) \quad (1)$$

where  $\Theta = (\theta_1, \theta_2, \theta_3)$  and  $K_4 = (\kappa_0, \kappa_1, \kappa_2, \kappa_3)$ . Below is the  $K_4$  representation with four diagonal matrices that their elements are in  $\{-1, 0, +1\}$  and  $\kappa_0$  is the identity,

$$\kappa_0 = \begin{pmatrix} +1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & +1 \end{pmatrix}, \kappa_1 = \begin{pmatrix} +1 & 0 & 0 & 0 \\ 0 & +1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}, \quad (2)$$

$$\kappa_2 = \begin{pmatrix} +1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & +1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}, \kappa_3 = \begin{pmatrix} +1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & +1 \end{pmatrix}, \quad (3)$$

where  $\det(\kappa_i) = 1$ ,  $\text{trace}(\kappa_i) = 0 (i \neq 0)$ , and  $\kappa_i$  is diagonal for all  $i \in \{0, 1, 2, 3\}$ . The absolute square of the triangle qubit  $T(\Theta, K_4)$  is

$$\begin{aligned} T(\Theta, K_4) T(\Theta, K_4)^* &= T(\Theta, K_4) T(\Theta, K_4)^\dagger \\ &= \frac{1}{2} (\kappa_0 + \kappa_1 \exp(+i2\pi\theta_1) + \kappa_2 \exp(+i2\pi\theta_2) + \kappa_3 \exp(+i2\pi\theta_3)) \\ &\quad \cdot \frac{1}{2} (\kappa_0 + \kappa_1 \exp(-i2\pi\theta_1) + \kappa_2 \exp(-i2\pi\theta_2) + \kappa_3 \exp(-i2\pi\theta_3)). \end{aligned}$$

Since  $K_4$  is an Abelian group with  $\kappa_0$  as identity,  $\kappa_i^2 = I_4$  for all  $i \in \{0, 1, 2, 3\}$ , and  $\kappa_1 \kappa_2 = \kappa_3, \kappa_1 \kappa_3 = \kappa_2, \kappa_2 \kappa_3 = \kappa_1$ , the

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absolute square of  $T(\Theta, K_4)$  is equal to

$$\begin{aligned} & \kappa_0 + \frac{1}{4} \{ \kappa_1 (\exp(i2\pi\theta_1) + \exp(-i2\pi\theta_1)) \\ & + \kappa_2 (\exp(i2\pi\theta_2) + \exp(-i2\pi\theta_2)) \\ & + \kappa_3 (\exp(i2\pi\theta_3) + \exp(-i2\pi\theta_3)) \\ & + \kappa_1 (\exp(i2\pi(\theta_2 - \theta_3)) + \exp(-i2\pi(\theta_2 - \theta_3))) \\ & + \kappa_2 (\exp(i2\pi(\theta_1 - \theta_3)) + \exp(-i2\pi(\theta_1 - \theta_3))) \\ & + \kappa_3 (\exp(i2\pi(\theta_1 - \theta_2)) + \exp(-i2\pi(\theta_1 - \theta_2))) \}. \end{aligned}$$

Furthermore, let  $\hat{\theta} = \frac{\theta_1 + \theta_2 + \theta_3}{2}$ , by Eluer's formula and cosine multiplication formula, the absolute square of  $T(\Theta, K_4)$  is equal to

$$\begin{aligned} & \kappa_0 + \kappa_1 \cos(\hat{\theta} - \theta_3) \cos(\hat{\theta} - \theta_2) + \kappa_2 \cos(\hat{\theta} - \theta_3) \cos(\hat{\theta} - \theta_1) \\ & + \kappa_3 \cos(\hat{\theta} - \theta_2) \cos(\hat{\theta} - \theta_1). \end{aligned}$$

Moreover, let  $\omega_1, \omega_2, \omega_3 \in \mathbb{R}$  such that  $\theta_1 = \omega_1 + \omega_3, \theta_2 = \omega_2 + \omega_3, \theta_3 = \omega_1 + \omega_2$ , we obtain  $\hat{\theta} = \omega_1 + \omega_2 + \omega_3, \omega_1 = \hat{\theta} - \theta_2, \omega_2 = \hat{\theta} - \theta_1$ , and  $\omega_3 = \hat{\theta} - \theta_3$ , the absolute square of  $U(\Theta, K_4)$  is equal to

$$\begin{aligned} & \kappa_0 + \kappa_1 \cos(2\pi\omega_1) \cos(2\pi\omega_3) + \kappa_2 \cos(2\pi\omega_2) \cos(2\pi\omega_3) \\ & + \kappa_3 \cos(2\pi\omega_1) \cos(2\pi\omega_2) \end{aligned}$$

Let  $(\lambda_1, \lambda_2, \lambda_3) = (\cos(2\pi\omega_1) \cos(2\pi\omega_3), \cos(2\pi\omega_2) \cos(2\pi\omega_3), \cos(2\pi\omega_1) \cos(2\pi\omega_2))$  and  $(q_0, q_1, q_2, q_3) = (\frac{1}{4}(1 + \lambda_1 + \lambda_2 + \lambda_3), \frac{1}{4}(1 + \lambda_1 - \lambda_2 - \lambda_3), \frac{1}{4}(1 - \lambda_1 + \lambda_2 - \lambda_3), \frac{1}{4}(1 - \lambda_1 - \lambda_2 + \lambda_3))$ , since  $T(\Theta; K_4)T(\Theta; K_4)^\dagger \geq 0$  which is equivalent to  $\text{diag}(q_0, q_1, q_2, q_3) \geq 0$  where the  $\frac{1}{4}(q_0, q_1, q_2, q_3)$  are eigenvalues of Choi's matrix<sup>7</sup> for density operator  $\rho = \frac{1}{2}(I_2 + \sum_{i=1}^3 \lambda_i \sigma_i)$ . In addition, since  $|\lambda_i| \leq 1$ , by Lemma 5.11<sup>4</sup> and Choi's theorem<sup>7</sup>, the diagonal matrix  $\text{diag}(1, \lambda_1, \lambda_2, \lambda_3)$  is a unital qubit channel.  $\square$

**Remark 1.** By the definition of  $(\lambda_1, \lambda_2, \lambda_3)$  above, there is  $\lambda_1 \lambda_2 \lambda_3 = (\cos(2\pi\omega_1) \cos(2\pi\omega_2) \cos(2\pi\omega_3))^2$ , thus  $\cos(2\pi\omega_1) \cos(2\pi\omega_2) \cos(2\pi\omega_3) = \pm \sqrt{\lambda_1 \lambda_2 \lambda_3}$ . We obtain  $\cos(2\pi\omega_1) = \pm \sqrt{\lambda_1 \lambda_3 / \lambda_2}$ ,  $\cos(2\pi\omega_2) = \pm \sqrt{\lambda_2 \lambda_3 / \lambda_1}$ , and  $\cos(2\pi\omega_3) = \pm \sqrt{\lambda_1 \lambda_2 / \lambda_3}$ . The map from  $\lambda_i$  to  $\omega_i$  is not bijective.

#### IV. ANTIPODAL POINT AND BELL'S INEQUALITY

**Theorem 2.** The antipodal point of a triangle qubit channel on Bloch sphere is non-violation Bell's inequality with cosine correlation functions.

*Proof.* Given a qubit channel with a point  $(\lambda_1, \lambda_2, \lambda_3)$  in Bloch sphere, its antipodal point is equal to  $(\lambda'_1, \lambda'_2, \lambda'_3) = (-\lambda_1, -\lambda_2, -\lambda_3)$  with the density operator  $\rho' = \frac{1}{2}(I_2 + \sum_{i=1}^3 \lambda'_i \sigma_i)$ . Choi's matrix of  $\rho'$  is

$$C_{\rho'} = \frac{1}{4} \begin{pmatrix} +1 - \lambda_3 & 0 & 0 & -\lambda_1 - \lambda_2 \\ 0 & +1 + \lambda_3 & -\lambda_1 + \lambda_2 & 0 \\ 0 & -\lambda_1 + \lambda_2 & +1 + \lambda_3 & 0 \\ -\lambda_1 - \lambda_2 & 0 & 0 & +1 - \lambda_3 \end{pmatrix}. \quad (4)$$

Let the eigenvalues of  $C_{\rho'}$  be  $(q'_0, q'_1, q'_2, q'_3)$ , we obtain  $q'_0 = \frac{1}{4}(1 - \lambda_1 - \lambda_2 - \lambda_3)$ ,  $q'_1 = \frac{1}{4}(1 - \lambda_1 + \lambda_2 + \lambda_3)$ ,  $q'_2 = \frac{1}{4}(1 + \lambda_1 - \lambda_2 + \lambda_3)$ ,  $q'_3 = \frac{1}{4}(1 + \lambda_1 + \lambda_2 - \lambda_3)$ . To the corresponding  $q_i$  in the proof of Theorem 1, for all  $i \in \{0, 1, 2, 3\}$  there are  $q_i + q'_i = \frac{1}{2}$  and  $0 \leq q_i \leq 1$  because  $|\lambda_i| \leq 1$  and  $q_i$  is an eigenvalue of Choi's matrix of the unital qubit channel. Thus, we have  $-\frac{1}{2} \leq q'_i \leq \frac{1}{2}$ .

To an unital qubit channel, there are  $q_1 = 1 + \lambda_1 - \lambda_2 - \lambda_3 \geq 0$  and  $q_2 = 1 - \lambda_1 + \lambda_2 - \lambda_3 \geq 0$ . That is,  $-(1 - \lambda_3) \leq \lambda_1 - \lambda_2 \leq 1 - \lambda_3$  which is  $|\lambda_1 - \lambda_2| \leq 1 - \lambda_3$ , one of a FAC inequality<sup>5</sup>. For the antipodal point since  $\lambda'_i = -\lambda_i$ , we have  $|\lambda'_1 - \lambda'_2| \leq 1 + \lambda'_3$ . When  $P_{xz}, P_{yz}, P_{xy}$  are correlation functions and  $\lambda'_1 = P_{xz}, \lambda'_2 = P_{yz}, \lambda'_3 = P_{xy}$ , we obtain Bell's inequality<sup>2</sup>:

$$|P_{xz} - P_{yz}| \leq 1 + P_{xy}. \quad (5)$$

Epecially, let the cosine correlation functions be  $P_{xz} = \cos(2\pi\omega_1) \cos(2\pi\omega_3)$ ,  $P_{yz} = \cos(2\pi\omega_2) \cos(2\pi\omega_3)$ , and  $P_{xy} = \cos(2\pi\omega_1) \cos(2\pi\omega_2)$ , we proved the theorem for the antipodal point.  $\square$

#### V. DISCUSSION

A special case in theorem 1 is  $(\omega_1, \omega_2, \omega_3) = (pt, qt, rt)$  where  $(p, q, r)$  are relatively prime and  $t \in \mathbb{R}$ , we obtain three cosine wave functions  $\cos(2\pi pt), \cos(2\pi qt)$ , and  $\cos(2\pi rt)$  with frequencies  $(p, q, r)$ . The triangle qubit matrix  $T(\Theta; K_4)$  becomes

$$T(z, p, q, r; K_4) = \frac{1}{2} (\kappa_0 + \kappa_1 z^{p+r} + \kappa_2 z^{q+r} + \kappa_3 z^{p+q}) \quad (6)$$

where  $z = \exp(i2\pi t)$ , which is a polynomial matrix over Klein four-group. In physics, we can apply three cosine waves with zero phases to construct a continuous triangle qubit channel.

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